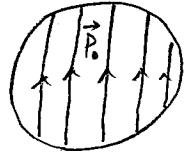


Date: 2014 Oct 27

- ① Let us consider a solid sphere, of radius R , with permanent polarization

$$\vec{P}(\vec{r}, t) = \vec{P}_0 \theta(R-r)$$



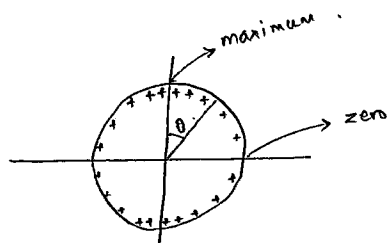
- ② Maxwell's equations are.

$$\vec{\nabla} \cdot (\epsilon_0 \vec{E}) = \rho - \underbrace{\vec{\nabla} \cdot \vec{P}}_{\rho_{\text{eff.}}}$$

$$\vec{\nabla} \times \vec{E} = 0$$

$$\begin{aligned} \textcircled{3} \quad \rho_{\text{eff}}(\vec{r}) &= - \vec{\nabla} \cdot \vec{P} \\ &= - \vec{\nabla} \cdot [\vec{P}_0 \theta(R-r)] \\ &= - \vec{P}_0 \cdot \vec{\nabla} \theta(R-r) \\ &= - [\vec{P}_0 \cdot (\vec{\nabla} r)] \frac{\partial}{\partial r} \theta(R-r) \\ &= (\vec{P}_0 \cdot \hat{r}) \delta(R-r) \\ &= P_0 \cos \theta \delta(R-r) \end{aligned}$$

$$\frac{\partial}{\partial x} \theta(R-r) = \frac{\partial r}{\partial x} \frac{\partial}{\partial r} \theta(R-r)$$

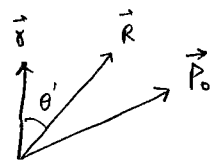


④ To find the electric potential we have.

$$-\nabla^2 \phi = \frac{1}{\epsilon_0} \rho_{\text{eff}}$$

$$\begin{aligned} \phi(\vec{r}) &= \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \frac{\rho_{\text{eff}}(\vec{r}')}{|\vec{r} - \vec{r}'|} \\ &= \frac{1}{4\pi\epsilon_0} \int_0^\infty r'^2 dr' \int_0^\pi \sin\theta' d\theta' \int_0^{2\pi} d\phi' \frac{\vec{P}_0 \cdot \{ \hat{i} \sin\theta' \cos\phi' + \hat{j} \sin\theta' \sin\phi' + \cos\theta' \hat{k} \}}{\sqrt{r^2 + r'^2 - 2\vec{r} \cdot \vec{r}'}} \\ &= \frac{R^2}{4\pi\epsilon_0} \int_0^\pi \sin\theta' d\theta' \int_0^{2\pi} d\phi' \frac{\vec{P}_0 \cdot \{ \hat{i} \sin\theta' \cos\phi' + \hat{j} \sin\theta' \sin\phi' + \cos\theta' \hat{k} \}}{\sqrt{r^2 + R^2 - 2\vec{r} \cdot \vec{R}}} \end{aligned}$$

⑤ Choose $\vec{r}$ to be along z-axis. If we do not make the choice, the integrals will require the use of properties of Legendre polynomials.



$$\vec{r} \cdot \vec{R} = rR \cos\theta'$$

⑥ This allows us to complete the ϕ' integral

$$\begin{aligned} \phi(\vec{r}) &= \frac{R^2}{4\pi\epsilon_0} \int_0^\pi \sin\theta' d\theta' \int_0^{2\pi} d\phi' \frac{\vec{P}_0 \cdot \{ \hat{i} \sin\theta' \cos\phi' + \hat{j} \sin\theta' \sin\phi' + \hat{k} \cos\theta' \}}{\sqrt{r^2 + R^2 - 2rR \cos\theta'}} \\ &= \frac{R^2}{4\pi\epsilon_0} (\vec{P}_0 \cdot \hat{k}) 2\pi \int_0^\pi \sin\theta' d\theta' \frac{\cos\theta'}{\sqrt{r^2 + R^2 - 2rR \cos\theta'}} \\ &= \frac{2\pi R^2}{4\pi\epsilon_0} (\vec{P}_0 \cdot \hat{k}) \int_{-1}^{+1} dt \frac{t}{\sqrt{r^2 + R^2 - 2rR t}} \end{aligned}$$

⑦

Homework :

Show that, (using substitution $r^2 + R^2 - 2rRt = y$.)

$$(i) \int_{-1}^{+1} dt \frac{1}{\sqrt{r^2 + R^2 - 2rRt}} =$$

$$(ii) \int_{-1}^{+1} dt \frac{t}{\sqrt{r^2 + R^2 - 2rRt}} = \begin{cases} \frac{2}{3} \frac{r}{R^2}, & r < R, \\ \frac{2}{3} \frac{R}{r^2}, & R < r \end{cases}$$

$$= \frac{2}{3} \frac{r}{r^2}$$

⑧ Using ⑦ in ⑥

$$\phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} (\vec{P}_0 \cdot \hat{k}) \begin{cases} \frac{4\pi}{3} r, & r < R, \\ \left(\frac{4\pi}{3} R^3\right) \frac{1}{r^2}, & R < r, \end{cases}$$

$$= \frac{1}{4\pi\epsilon_0} (\vec{P}_0 \cdot \hat{r}) \begin{cases} \left(\frac{4\pi}{3} r^3\right) \frac{1}{r^2}, & r < R, \\ \left(\frac{4\pi}{3} R^3\right) \frac{1}{r^2}, & R < r. \end{cases}$$

(replaced $\hat{k}$ with $\hat{r}$, as per our choice in ⑤.)

④ The electric field is obtained using

$$\vec{E} = -\vec{\nabla} \phi$$

$$= \frac{1}{4\pi\epsilon_0} \begin{cases} \frac{4\pi}{3} (-1) \vec{\nabla} [\vec{P}_0 \cdot \vec{r}], & r < R, \\ \frac{4\pi}{3} (-1) R^3 \vec{\nabla} \left(\frac{\vec{P}_0 \cdot \vec{r}}{r^3} \right), & R < r, \end{cases}$$

$$= \frac{1}{4\pi\epsilon_0} \begin{cases} -\frac{4\pi}{3} \vec{P}_0, & r < R, \\ \left(\frac{4\pi}{3} R^3 \right) \frac{1}{r^3} [3(\vec{P}_0 \cdot \hat{r}) \hat{r} - \vec{P}_0], & R < r. \end{cases}$$

Thus, outside the sphere the electric field is that of a point dipole. Inside the sphere it is a constant. That is, the charge accumulated on the surface is such that it produces uniform $\vec{E}$, which simulates a dipole.

