Homework No. 10 (Spring 2017)

PHYS 510: Classical Mechanics

Due date: 2017 May 4 (Thursday) 4.30pm

1. (20 points.) The Lorentz factor

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}}, \qquad \beta = \frac{v}{c}.$$
(1)

- (a) Evaluate γ for v = 30 m/s (~ 70 miles/hour).
- (b) Evaluate γ for v = 3c/5.
- 2. (20 points.) Lorentz transformation describing a boost in the x-direction, y-direction, and z-direction, are

$$L_{1} = \begin{pmatrix} \gamma_{1} & -\beta_{1}\gamma_{1} & 0 & 0 \\ -\beta_{1}\gamma_{1} & \gamma_{1} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad L_{2} = \begin{pmatrix} \gamma_{2} & 0 & -\beta_{2}\gamma_{2} & 0 \\ 0 & 1 & 0 & 0 \\ -\beta_{2}\gamma_{2} & 0 & \gamma_{2} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad L_{3} = \begin{pmatrix} \gamma_{3} & 0 & 0 & -\beta_{3}\gamma_{3} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\beta_{3}\gamma_{3} & 0 & 0 & \gamma_{3} \end{pmatrix},$$

$$(2)$$

respectively. Transformation describing a rotation about the x-axis, y-axis, and z-axis, are

$$R_{1} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \cos \omega_{1} & \sin \omega_{1} \\ 0 & 0 & -\sin \omega_{1} & \cos \omega_{1} \end{pmatrix}, \quad R_{2} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \omega_{2} & 0 & -\sin \omega_{2} \\ 0 & 0 & 1 & 0 \\ 0 & \sin \omega_{2} & 0 & \cos \omega_{2} \end{pmatrix}, \quad R_{3} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \omega_{3} & \sin \omega_{3} & 0 \\ 0 & -\sin \omega_{3} & \cos \omega_{3} & 0 \\ 0 & 0 & 0 & 1 \\ (3) \end{pmatrix},$$

respectively. For infinitesimal transformations, $\beta_i = \delta \beta_i$ and $\omega_i = \delta \omega_i$ use the approximations

$$\gamma_i \sim 1, \qquad \cos \omega_i \sim 1, \qquad \sin \omega_i \sim \delta \omega_i,$$
(4)

to identify the generator for boosts \mathbf{N} , and the generator for rotations the angular momentum \mathbf{J} ,

$$\mathbf{L} = \mathbf{1} + \delta \boldsymbol{\beta} \cdot \mathbf{N} \quad \text{and} \quad \mathbf{R} = \mathbf{1} + \delta \boldsymbol{\omega} \cdot \mathbf{J}, \tag{5}$$

respectively. Then derive

$$[N_1, N_2] = N_1 N_2 - N_2 N_1 = J_3.$$
(6)

This states that boosts in perpendicular direction leads to rotation. (To gain insight of the statement, calculate $[J_1, J_2]$ and interpret the result.)

3. (60 points.) The Poincaré formula for the addition of (parallel) velocities is

$$v = \frac{v_a + v_b}{1 + \frac{v_a v_b}{c^2}},\tag{7}$$

where v_a and v_b are velocities and c is speed of light in vacuum. Jerzy Kocik, from the department of Mathematics in SIUC, has invented a geometric diagram that allows one to visualize the Poincaré formula. (Refer [1].) An interactive applet for exploring velocity addition is available at Kocik's web page [2]. (For the following assume that the Poincaré formula holds for all speeds, subluminal $(v_i < c)$, superluminal $(v_i > c)$, and speed of light.)

- (a) Analyse what is obtained if you add two subluminal speeds?
- (b) Analyse what is obtained if you add a subluminal speed to speed of light?
- (c) Analyse what is obtained if you add a subluminal speed to a superluminal speed?
- (d) Analyse what is obtained if you add speed of light to another speed of light?
- (e) Analyse what is obtained if you add a superluminal speed to speed of light?
- (f) Analyse what is obtained if you add two superluminal speeds?
- 4. (**30 points.**) Let

$$\tanh \theta = \beta, \tag{8}$$

where $\beta = v/c$. Addition of (parallel) velocities in terms of the parameter θ obeys the arithmatic addition

$$\theta = \theta_a + \theta_b. \tag{9}$$

- (a) Invert the expression in Eq. (8) to find the explicit form of θ in terms of β as a logarithm.
- (b) Show that Eq. (9) leads to the relation

$$\left(\frac{1+\beta}{1-\beta}\right) = \left(\frac{1+\beta_a}{1-\beta_a}\right) \left(\frac{1+\beta_b}{1-\beta_b}\right). \tag{10}$$

(c) Using Eq. (10) derive the Poincaré formula for the addition of (parallel) velocities.

5. (**20 points.**) (Refer Hughston and Tod's book.) Prove that

- (a) if p_{μ} is a time-like vector and $p^{\mu}s_{\mu} = 0$ then s^{μ} is necessarily space-like.
- (b) if p_{μ} and q^{μ} are both time-like vectors and $p^{\mu}q_{\mu} > 0$ then either both are futurepointing or both are past-pointing.
- 6. (100 points.) Eigenvalues of the energy momentum tensor. (We choose c = 1, which is easily undone by replacing $\mathbf{E} \to \mathbf{E}/c$ everywhere.)

(a) Using

$$F_{\mu\nu} = \begin{pmatrix} 0 & -E_1 & -E_2 & -E_3 \\ E_1 & 0 & B_3 & -B_2 \\ E_2 & -B_3 & 0 & B_1 \\ E_3 & B_2 & -B_1 & 0 \end{pmatrix}$$
(11)

and

$$\tilde{F}_{\mu\nu} = \frac{1}{2} \varepsilon_{\mu\nu\alpha\beta} F^{\alpha\beta} \tag{12}$$

evaluate the following:

i. $F^{\mu\lambda}F_{\lambda\nu}$

ii. $\tilde{F}^{\mu\lambda}\tilde{F}_{\lambda\nu}$

iii. Then, derive

$$F^{\mu\lambda}\tilde{F}_{\lambda\nu} = \delta^{\mu}{}_{\nu}\mathbf{E}\cdot\mathbf{B},\tag{13a}$$

$$\tilde{F}^{\mu\lambda}\tilde{F}_{\lambda\nu} - F^{\mu\lambda}F_{\lambda\nu} = \delta^{\mu}{}_{\nu}(B^2 - E^2).$$
(13b)

(b) Define

$$2f = (B^2 - E^2)$$
 and $g = \mathbf{E} \cdot \mathbf{B}$. (14)

Thus, construct matrix (or dyadic) equations

$$\mathbf{F} \cdot \ddot{\mathbf{F}} = g\mathbf{1},\tag{15a}$$

$$\tilde{\mathbf{F}} \cdot \tilde{\mathbf{F}} - \mathbf{F} \cdot \mathbf{F} = 2f\mathbf{1},\tag{15b}$$

in terms of matrices (or dyadics) **F** and **F**.

(c) Show that the eigenvalues λ of the field tensor **F** satisfy the quartic equation

$$\lambda^4 + 2f\lambda^2 - g^2 = 0. (16)$$

(d) Evaluate the eigenvalues to be $\pm \lambda_1$ and $\pm \lambda_2$ where

$$\lambda_1 = \sqrt{-f - \sqrt{f^2 + g^2}} = \frac{i}{\sqrt{2}} \left[\sqrt{f + ig} + \sqrt{f - ig} \right], \tag{17}$$

$$\lambda_2 = \sqrt{-f + \sqrt{f^2 + g^2}} = \frac{i}{\sqrt{2}} \left[\sqrt{f + ig} - \sqrt{f - ig} \right]. \tag{18}$$

(e) Show that

i. if $B^2 - E^2 = 0$, then the eigenvalues are $\pm \sqrt{g}$ and $\pm i\sqrt{g}$.

- ii. if $\mathbf{B} \cdot \mathbf{E} = 0$, then the eigenvalues are 0, 0, and $\pm \sqrt{2f}$.
- (f) Prove the following:
 - i. There is no Lorentz transformation connecting two reference frames such that the field is purely magnetic in origin in one and purely electric in origin in the other.

- ii. If $B^2 E^2 > 0$ in a frame, then there exists a frame in which the field is purely magnetic.
- iii. If $B^2-E^2<0$ in a frame, then there exists a frame in which the field is purely electric.
- iv. If $B^2 E^2 = 0$ in a frame, then there exists a frame in which
 - **B** is perpendicular to **E**, if $\mathbf{B} \cdot \mathbf{E} = 0$.
 - **B** is parallel to **E**, if $\mathbf{B} \cdot \mathbf{E} > 0$.
 - **B** is antiparallel to **E**, if $\mathbf{B} \cdot \mathbf{E} < 0$.

References

- [1] J. Kocik. Geometric diagram for relativistic addition of velocities. Am. J. Phys., 80:737–739, August 2012.
- [2] J. Kocik. An interactive applet for exploring relativistic velocity addition. http://lagrange.math.siu.edu/Kocik/relativity/diagram.html.