

# Homework No. 09 (Fall 2022)

## PHYS 500A: MATHEMATICAL METHODS

*School of Physics and Applied Physics, Southern Illinois University–Carbondale*

Due date: Monday, 2022 Nov 28, 4.30pm

1. (10 points.) Integral representations for the modified Bessel functions,  $I_m(t)$  and  $K_m(t)$ , for integer  $m$  and  $0 \leq t < \infty$  are

$$K_m(t) = \int_0^\infty d\theta \cosh m\theta e^{-t \cosh \theta}, \quad (1a)$$

$$I_m(t) = \int_0^\pi \frac{d\phi}{\pi} \cos m\phi e^{t \cos \phi}. \quad (1b)$$

- (a) Using Mathematica (or your favourite graphing tool) plot  $K_0(t), K_1(t), K_2(t)$  and  $I_0(t), I_1(t), I_2(t)$  on the same plot. (Please do not submit hand sketched plots.)
- (b) Refer Chapter 10 of Digital Library of Mathematical Functions,

<https://dlmf.nist.gov/10>

for a comprehensive resource.

2. (10 points.) Show that the integral representations for the modified Bessel functions,  $I_m(t)$  and  $K_m(t)$ , for integer  $m$  and  $0 \leq t < \infty$ ,

$$K_m(t) = \int_0^\infty d\theta \cosh m\theta e^{-t \cosh \theta}, \quad (2a)$$

$$I_m(t) = \int_0^\pi \frac{d\phi}{\pi} \cos m\phi e^{t \cos \phi}. \quad (2b)$$

satisfies the differential equation for modified Bessel functions,

$$\left[ -\frac{1}{t} \frac{d}{dt} t \frac{d}{dt} + \frac{m^2}{t^2} + 1 \right] \begin{Bmatrix} I_m(t) \\ K_m(t) \end{Bmatrix} = 0. \quad (3)$$

Hint: Integrate by parts, after identifying

$$(t \cosh \theta - t^2 \sinh^2 \theta) e^{-t \cosh \theta} = -\frac{d^2}{d\theta^2} e^{-t \cosh \theta}, \quad (4a)$$

$$(t \cos \phi - t^2 \sin^2 \phi) e^{t \cos \phi} = -\frac{d^2}{d\phi^2} e^{t \cos \phi}. \quad (4b)$$

3. (20 points.) The modified Bessel functions,  $I_m(t)$  and  $K_m(t)$ , satisfy the differential equation

$$\left[ -\frac{1}{t} \frac{d}{dt} t \frac{d}{dt} + \frac{m^2}{t^2} + 1 \right] \begin{Bmatrix} I_m(t) \\ K_m(t) \end{Bmatrix} = 0. \quad (5)$$

Derive the identity, for the Wronskian, (upto a constant  $C$ )

$$I_m(t)K'_m(t) - K_m(t)I'_m(t) = -\frac{C}{t}, \quad (6)$$

where

$$I'_m(t) \equiv \frac{d}{dt} I_m(t) \quad \text{and} \quad K'_m(t) \equiv \frac{d}{dt} K_m(t). \quad (7)$$

Further, determine the value of the constant  $C$  on the right hand side of Eq. (6) using the asymptotic forms for the modified Bessel functions:

$$I_m(t) \xrightarrow{t \gg 1} \frac{1}{\sqrt{2\pi}} \frac{e^t}{\sqrt{t}}, \quad (8)$$

$$K_m(t) \xrightarrow{t \gg 1} \sqrt{\frac{\pi}{2}} \frac{e^{-t}}{\sqrt{t}}. \quad (9)$$

4. (20 points.) The cylindrical free Green's function satisfies

$$\left[ -\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} + \frac{m^2}{\rho^2} + k_z^2 \right] g_m(\rho, \rho'; k_z) = \frac{\delta(\rho - \rho')}{\rho}. \quad (10)$$

Integrate Eq. (10) around  $\rho = \rho'$  to derive the continuity conditions:

$$g_m(\rho, \rho'; k_z) \Big|_{\rho=\rho'-\delta}^{\rho=\rho'+\delta} = 0, \quad (11a)$$

$$\rho \frac{\partial}{\partial \rho} g_m(\rho, \rho'; k_z) \Big|_{\rho=\rho'-\delta}^{\rho=\rho'+\delta} = -1. \quad (11b)$$

Let us further require that

$$g_m(0, \rho'; k_z) \text{ is finite,} \quad (12a)$$

$$g_m(\infty, \rho'; k_z) = 0. \quad (12b)$$

Recall the Wronskian

$$I_m(t)K'_m(t) - I'_m(t)K_m(t) = -\frac{1}{t}. \quad (13)$$

Construct the solution to have the form

$$g_m(\rho, \rho') = \begin{cases} A I_m(k_z \rho) + B K_m(k_z \rho), & 0 \leq \rho < \rho', \\ C I_m(k_z \rho) + D K_m(k_z \rho), & \rho' < \rho < \infty. \end{cases} \quad (14)$$

Derive the solution

$$g_m(\rho, \rho') = I_m(k_z \rho_{<}) K_m(k_z \rho_{>}), \quad (15)$$

where  $\rho_{<} = \text{Minimum}(\rho, \rho')$  and  $\rho_{>} = \text{Maximum}(\rho, \rho')$ .

5. (20 points.) The cylindrical Green's function satisfies

$$\left[ -\frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} + \frac{m^2}{\rho^2} + k_z^2 \right] g_m(\rho, \rho'; k_z) = \frac{\delta(\rho - \rho')}{\rho}. \quad (16)$$

Inside a perfectly conducting cylinder we require

$$g_m(0, \rho'; k_z) \text{ is finite,} \quad (17a)$$

$$g_m(a, \rho'; k_z) = 0. \quad (17b)$$

This shields all of (electrostatics related) physics between the plates from outside. Thus, we have

$$0 \leq \rho \leq a, \quad (18a)$$

$$0 \leq \rho' \leq a, \quad (18b)$$

Integrate Eq. (16) around  $\rho = \rho'$  to derive the continuity conditions:

$$g_m(\rho, \rho'; k_z) \Big|_{\rho=\rho'-\delta}^{\rho=\rho'+\delta} = 0, \quad (19a)$$

$$\rho \frac{\partial}{\partial \rho} g_m(\rho, \rho'; k_z) \Big|_{\rho=\rho'-\delta}^{\rho=\rho'+\delta} = -1. \quad (19b)$$

Recall the Wronskian

$$I_m(t)K'_m(t) - I'_m(t)K_m(t) = -\frac{1}{t}. \quad (20)$$

Construct the solution to have the form

$$g_m(\rho, \rho') = \begin{cases} A I_m(k_z \rho) + B K_m(k_z \rho), & 0 \leq \rho < \rho', \\ C I_m(k_z \rho) + D K_m(k_z \rho), & \rho' < \rho < a. \end{cases} \quad (21)$$

Derive the solution

$$g_m(\rho, \rho') = I_m(k_z \rho_{<}) K_m(k_z \rho_{>}) - \frac{K_m(k_z a)}{I_m(k_z a)} I_m(k_z \rho) I_m(k_z \rho'). \quad (22)$$