Homework No. 03 (Fall 2025)

PHYS 500A: MATHEMATICAL METHODS

School of Physics and Applied Physics, Southern Illinois University-Carbondale

Due date: Monday, 2025 Sep 8, 4.30pm

- 0. Problems 1 and 3 are for practice. Problem 2 is for submission.
- 1. (**Example.**) Let \mathbf{r} represent the position vector, x^i the components of the position vector in rectangular coordinates, and u^i the components of the position vector in cylindrical polar coordinates. In particular, we have

$$x^{1} = x = \rho \cos \phi,$$
 $u^{1} = \rho = \sqrt{x^{2} + y^{2}},$ (1a)

$$x^{2} = y = \rho \sin \phi,$$
 $u^{2} = \phi = \tan^{-1} \frac{y}{x},$ (1b)

$$x^3 = z = z,$$
 $u^3 = z = z.$ (1c)

Let us define the unit vectors

$$\hat{\boldsymbol{\rho}} = \cos\phi \,\hat{\mathbf{i}} + \sin\phi \,\hat{\mathbf{j}} + 0 \,\hat{\mathbf{k}},\tag{2a}$$

$$\hat{\boldsymbol{\phi}} = -\sin\phi\,\hat{\mathbf{i}} + \cos\phi\,\hat{\mathbf{j}} + 0\,\hat{\mathbf{k}},\tag{2b}$$

$$\hat{\mathbf{z}} = 0\,\hat{\mathbf{i}} + 0\,\hat{\mathbf{j}} + \hat{\mathbf{k}},\tag{2c}$$

where $\hat{\mathbf{i}}$, $\hat{\mathbf{j}}$, and $\hat{\mathbf{k}}$ are basis vectors in rectangular coordinate system. We will also use the notation $\hat{\mathbf{i}} = \hat{\mathbf{x}}^1 = \hat{\mathbf{x}}_1$, $\hat{\mathbf{j}} = \hat{\mathbf{x}}^2 = \hat{\mathbf{x}}_2$, $\hat{\mathbf{k}} = \hat{\mathbf{x}}^3 = \hat{\mathbf{x}}_3$, to represent these vectors.

(a) Basis vectors:

$$\mathbf{e}_i = \frac{\partial \mathbf{r}}{\partial u^i}.\tag{3}$$

Show that

$$\mathbf{e}_1 = \hat{\boldsymbol{\rho}}, \qquad \mathbf{e}_2 = \rho \hat{\boldsymbol{\phi}}, \qquad \mathbf{e}_3 = \hat{\mathbf{z}}.$$
 (4)

(b) Reciprocal basis vectors:

$$\mathbf{e}^i = \mathbf{\nabla} u^i. \tag{5}$$

Show that

$$\mathbf{e}^1 = \hat{\boldsymbol{\rho}}, \qquad \mathbf{e}^2 = \frac{\hat{\boldsymbol{\phi}}}{\rho}, \qquad \mathbf{e}^3 = \hat{\mathbf{z}}.$$
 (6)

(c) Orthonormality: Show that

$$\mathbf{e}^i \cdot \mathbf{e}_j = \delta^i_j. \tag{7}$$

(d) Metric tensor: A line element is defined as

$$d\mathbf{r} = dx^i \hat{\mathbf{x}}_i = du^i \mathbf{e}_i. \tag{8}$$

Show that

$$d\mathbf{r} \cdot d\mathbf{r} = du^i du^j g_{ii},\tag{9}$$

where the metric tensor g_{ij} is defined as

$$g_{ij} = \mathbf{e}_i \cdot \mathbf{e}_j. \tag{10}$$

Evaluate all the components of g_{ij} .

(e) Completeness relation: Starting from

$$\nabla \mathbf{r} = \mathbf{1} \tag{11}$$

derive the completeness relation

$$\mathbf{e}^i \, \mathbf{e}_i = \mathbf{1}. \tag{12}$$

Express the completeness relation in cylindrical polar coordinates in terms of $\hat{\rho}$, $\hat{\phi}$, and $\hat{\mathbf{z}}$.

- 2. (20 points.) Transformation of basis vectors.
 - (a) Let us consider a set of basis vectors \mathbf{e}_i , where i = 1, 2, 3, and the associated reciprocal basis vectors \mathbf{e}^i that satisfy the completeness relation

$$1 = e^{i} e_{i} = e^{1} e_{1} + e^{2} e_{2} + e^{3} e_{3}.$$
(13)

It is a complete set because an arbitrary vector \mathbf{A} can be expressed in terms of its projections along the basis vectors in the following way,

$$\mathbf{A} = \mathbf{A} \cdot \mathbf{1} = \mathbf{A} \cdot (\mathbf{e}^i \, \mathbf{e}_i) = (\mathbf{A} \cdot \mathbf{e}^i) \, \mathbf{e}_i = A^i \, \mathbf{e}_i, \tag{14}$$

where we recognized and defined the projections of vector \mathbf{A} along the direction of basis vectors as the components

$$A^i = (\mathbf{A} \cdot \mathbf{e}^i). \tag{15}$$

Similarly, multiplying by the identity dyadic on the left gives

$$\mathbf{A} = \mathbf{1} \cdot \mathbf{A} = (\mathbf{e}^i \, \mathbf{e}_i) \cdot \mathbf{A} = \mathbf{e}^i \, (\mathbf{e}_i \cdot \mathbf{A}) = \mathbf{e}^i \, A_i, \tag{16}$$

where now the projections of vector \mathbf{A} in the direction of the reciprocal basis vectors are the components

$$A_i = (\mathbf{e}_i \cdot \mathbf{A}). \tag{17}$$

For consistency we require the equality

$$(\mathbf{A} \cdot \mathbf{e}^i) \, \mathbf{e}_i = \mathbf{e}^i \, (\mathbf{e}_i \cdot \mathbf{A}). \tag{18}$$

Thus, derive

$$A^j = g^{ji}A_i, (19a)$$

$$A_j = g_{ji}A^i, (19b)$$

where the metric tensors are defined as where the metric tensors are defined as

$$g^{ji} = \mathbf{e}^j \cdot \mathbf{e}^i, \tag{20a}$$

$$g_{ji} = \mathbf{e}_j \cdot \mathbf{e}_i. \tag{20b}$$

(b) For another set of basis vectors \mathbf{g}_i and the associated reciprocal basis vectors \mathbf{g}^i that satisfy the completeness relation

$$\mathbf{1} = \mathbf{g}^i \, \mathbf{g}_i \tag{21}$$

we can write

$$\mathbf{A} = (\mathbf{A} \cdot \mathbf{g}^i) \, \mathbf{g}_i = \bar{A}^i \, \mathbf{g}_i, \tag{22}$$

where the components \bar{A}^i are in general different from A^i , and

$$\mathbf{A} = \mathbf{g}^i \left(\mathbf{g}_i \cdot \mathbf{A} \right) = \mathbf{g}^i \, \bar{A}_i. \tag{23}$$

For consistency we require the equality

$$(\mathbf{A} \cdot \mathbf{g}^i) \, \mathbf{g}_i = \mathbf{g}^i \, (\mathbf{g}_i \cdot \mathbf{A}). \tag{24}$$

Thus, derive

$$\bar{A}^j = \bar{g}^{ji}\bar{A}_i, \tag{25a}$$

$$\bar{A}_j = \bar{g}_{ji}\bar{A}^i, \tag{25b}$$

where the metric tensors are defined as

$$\bar{g}^{ji} = \mathbf{g}^j \cdot \mathbf{g}^i, \tag{26a}$$

$$\bar{g}_{ji} = \mathbf{g}_j \cdot \mathbf{g}_i. \tag{26b}$$

(c) For consistency between the two independent basis vector representations we require the equality

$$(\mathbf{A} \cdot \mathbf{e}^i) \, \mathbf{e}_i = (\mathbf{A} \cdot \mathbf{g}^i) \, \mathbf{g}_i. \tag{27}$$

Taking the dot product on the right with e^{j} and using orthogonality relation

$$\mathbf{e}_i \cdot \mathbf{e}^i = \delta_i^i \tag{28}$$

we obtain

$$(\mathbf{A} \cdot \mathbf{e}^j) = (\mathbf{A} \cdot \mathbf{g}^i) [\mathbf{g}_i \cdot \mathbf{e}^j]. \tag{29}$$

Similarly, taking the dot product on the right with \mathbf{g}^{j} we obtain

$$(\mathbf{A} \cdot \mathbf{e}^i) \left[\mathbf{e}_i \cdot \mathbf{g}^j \right] = (\mathbf{A} \cdot \mathbf{g}^j). \tag{30}$$

In terms of the transformation matrices connecting the two basis vectors,

$$S_i{}^j = [\mathbf{g}_i \cdot \mathbf{e}^j] \tag{31}$$

and

$$R_i{}^j = [\mathbf{e}_i \cdot \mathbf{g}^j], \tag{32}$$

we can derive the transformation of the components

$$A^j = \bar{A}^i S_i^{\ j},\tag{33a}$$

$$A^i R_i{}^j = \bar{A}^j. (33b)$$

Further, derive

$$A_j = R_j^{\ i} \bar{A}_i, \tag{34a}$$

$$S_j^i A_i = \bar{A}^j. \tag{34b}$$

Show that

$$R_i{}^j S_j{}^k = [\mathbf{e}_i \cdot \mathbf{g}^j][\mathbf{g}_j \cdot \mathbf{e}^k] = \mathbf{e}_i \cdot (\mathbf{g}^j \, \mathbf{g}_j) \cdot \mathbf{e}^k = \mathbf{e}_i \cdot \mathbf{1} \cdot \mathbf{e}^k = \delta_i^k. \tag{35}$$

- (d) i. Find the transformation matrices S and T between cylindrical polar coordinates and rectangular coordinates. Verify that ST = 1.
 - ii. Find the transformation matrices S and T between cylindrical polar coordinates and spherical polar coordinates. Verify that ST = 1.
- 3. (**Example.**) Let **r** represent a position vector in three dimensional space. Let x^i be the components of the position vector in rectangular coordinates, which can be interpreted as surfaces of constant x^i . Let us coordinatize the space using the planes, labeled using β ,

$$y = mx + \beta \tag{36}$$

where m is fixed, instead of planes with constant y. The other two sets of planes of constant x and constant z are the same. See Fig. 1. Let u^i be the components of the position vector in this new coordinatization of space. In particular, we have

$$x^1 = x = \alpha, u^1 = \alpha = x, (37a)$$

$$x^{2} = y = mx + \beta,$$
 $u^{2} = \beta = y - mx,$ (37b)

$$x^3 = z = \gamma, u^3 = \gamma = z. (37c)$$

The basis vectors $\hat{\mathbf{i}}$, $\hat{\mathbf{j}}$, and $\hat{\mathbf{k}}$ in rectangular coordinate system will be represented as $\hat{\mathbf{i}} = \hat{\mathbf{x}}^1 = \hat{\mathbf{x}}_1$, $\hat{\mathbf{j}} = \hat{\mathbf{x}}^2 = \hat{\mathbf{x}}_2$, $\hat{\mathbf{k}} = \hat{\mathbf{x}}^3 = \hat{\mathbf{x}}_3$, if necessary.

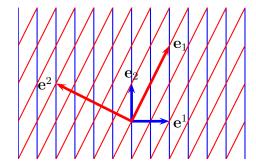


Figure 1: Basis vectors \mathbf{e}_i and reciprocal basis vectors \mathbf{e}^i .

(a) Basis vectors:

$$\mathbf{e}_{i} = \frac{\partial \mathbf{r}}{\partial u^{i}}.\tag{38}$$

Show that

$$\mathbf{e}_1 = \hat{\mathbf{i}} + m\,\hat{\mathbf{j}}, \qquad \mathbf{e}_2 = \hat{\mathbf{j}}, \qquad \mathbf{e}_3 = \hat{\mathbf{k}}.$$
 (39)

(b) Reciprocal basis vectors:

$$\mathbf{e}^i = \mathbf{\nabla} u^i. \tag{40}$$

Show that

$$\mathbf{e}^1 = \hat{\mathbf{i}}, \qquad \mathbf{e}^2 = -m\,\hat{\mathbf{i}} + \hat{\mathbf{j}}, \qquad \mathbf{e}^3 = \hat{\mathbf{k}}.$$
 (41)

Verify the relations

$$\mathbf{e}^{1} = \frac{\mathbf{e}_{2} \times \mathbf{e}_{3}}{(\mathbf{e}_{2} \times \mathbf{e}_{3}) \cdot \mathbf{e}_{1}}, \qquad \mathbf{e}^{2} = \frac{\mathbf{e}_{3} \times \mathbf{e}_{1}}{(\mathbf{e}_{3} \times \mathbf{e}_{1}) \cdot \mathbf{e}_{2}}, \qquad \mathbf{e}^{3} = \frac{\mathbf{e}_{1} \times \mathbf{e}_{2}}{(\mathbf{e}_{1} \times \mathbf{e}_{2}) \cdot \mathbf{e}_{3}}. \tag{42}$$

(c) Orthonormality: Show that

$$\mathbf{e}^i \cdot \mathbf{e}_j = \delta^i_j. \tag{43}$$

That is,

$$\mathbf{e}^1 \cdot \mathbf{e}_1 = 1,$$
 $\mathbf{e}^1 \cdot \mathbf{e}_2 = 0,$ $\mathbf{e}^1 \cdot \mathbf{e}_3 = 0,$ (44a)

$$\mathbf{e}^{2} \cdot \mathbf{e}_{1} = 0,$$
 $\mathbf{e}^{2} \cdot \mathbf{e}_{2} = 1,$ $\mathbf{e}^{2} \cdot \mathbf{e}_{3} = 0,$ (44b)
 $\mathbf{e}^{3} \cdot \mathbf{e}_{1} = 0,$ $\mathbf{e}^{3} \cdot \mathbf{e}_{2} = 0,$ $\mathbf{e}^{3} \cdot \mathbf{e}_{3} = 1.$ (44c)

$$e^{3} \cdot e_{1} = 0,$$
 $e^{3} \cdot e_{2} = 0,$ $e^{3} \cdot e_{3} = 1.$ (44c)

(d) Metric tensor: The metric tensor g_{ij} is defined as

$$g_{ij} = \mathbf{e}_i \cdot \mathbf{e}_j. \tag{45}$$

Evaluate all the components of g_{ij} . That is,

$$g_{11} = \mathbf{e}_1 \cdot \mathbf{e}_1 = 1 + m^2, \qquad g_{12} = \mathbf{e}_1 \cdot \mathbf{e}_2 = m, \qquad g_{13} = \mathbf{e}_1 \cdot \mathbf{e}_3 = 0,$$
 (46a)

$$g_{21} = \mathbf{e}_2 \cdot \mathbf{e}_1 = m,$$
 $g_{22} = \mathbf{e}_2 \cdot \mathbf{e}_2 = 1,$ $g_{23} = \mathbf{e}_2 \cdot \mathbf{e}_3 = 0,$ (46b)
 $g_{31} = \mathbf{e}_3 \cdot \mathbf{e}_1 = 0,$ $g_{32} = \mathbf{e}_3 \cdot \mathbf{e}_2 = 0,$ $g_{33} = \mathbf{e}_3 \cdot \mathbf{e}_3 = 1.$ (46c)

$$g_{31} = \mathbf{e}_3 \cdot \mathbf{e}_1 = 0,$$
 $g_{32} = \mathbf{e}_3 \cdot \mathbf{e}_2 = 0,$ $g_{33} = \mathbf{e}_3 \cdot \mathbf{e}_3 = 1.$ (46c)

Similarly evaluate the components of

$$g^{ij} = \mathbf{e}^i \cdot \mathbf{e}^j. \tag{47}$$

That is,

$$g^{11} = \mathbf{e}^{1} \cdot \mathbf{e}^{1} = 1, \qquad g^{12} = \mathbf{e}^{1} \cdot \mathbf{e}^{2} = -m, \qquad g^{13} = \mathbf{e}^{1} \cdot \mathbf{e}^{3} = 0, \qquad (48a)$$

$$g^{21} = \mathbf{e}^{2} \cdot \mathbf{e}^{1} = -m, \qquad g^{22} = \mathbf{e}^{2} \cdot \mathbf{e}^{2} = 1 + m^{2}, \qquad g^{23} = \mathbf{e}^{2} \cdot \mathbf{e}^{3} = 0, \qquad (48b)$$

$$g^{31} = \mathbf{e}^{3} \cdot \mathbf{e}^{1} = 0, \qquad g^{32} = \mathbf{e}^{3} \cdot \mathbf{e}^{2} = 0, \qquad g^{33} = \mathbf{e}^{3} \cdot \mathbf{e}^{3} = 1. \qquad (48c)$$

$$g^{21} = \mathbf{e}^2 \cdot \mathbf{e}^1 = -m, \qquad g^{22} = \mathbf{e}^2 \cdot \mathbf{e}^2 = 1 + m^2, \qquad g^{23} = \mathbf{e}^2 \cdot \mathbf{e}^3 = 0,$$
 (48b)

$$g^{31} = \mathbf{e}^3 \cdot \mathbf{e}^1 = 0,$$
 $g^{32} = \mathbf{e}^3 \cdot \mathbf{e}^2 = 0,$ $g^{33} = \mathbf{e}^3 \cdot \mathbf{e}^3 = 1.$ (48c)

Verify that $g^{ij}g_{jk} = \delta^i_k$.

(e) Completeness relation: Verify the completeness relation

$$\mathbf{e}^i \mathbf{e}_i = \mathbf{1} \tag{49}$$

by evaluating

$$e^1e_1 + e^2e_2 + e^3e_3.$$
 (50)

(f) Given a vector

$$\mathbf{A} = a\,\hat{\mathbf{i}} + b\,\hat{\mathbf{j}} + c\,\hat{\mathbf{k}} \tag{51}$$

in rectangular coordinates, find the components of the vector \mathbf{A} in the basis of \mathbf{e}_i . That is, find the components A^i in

$$\mathbf{A} = A^1 \,\mathbf{e}_1 + A^2 \,\mathbf{e}_2 + A^3 \,\mathbf{e}_3. \tag{52}$$